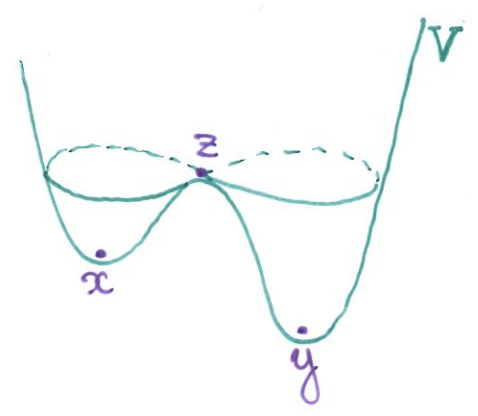


# Metastability in stochastic PDEs

Based on: NB & Barbara Gentz, EJP 18 (24): 1-58 (2013)  
 NB, Giacomo Di Gesù & Hendrik Weber, EJP 22 (41): 1-27 (2017)

## 1. Metastability in gradient SDEs

$$dx_t = -\nabla V(x_t) dt + \sqrt{2\varepsilon} dW_t$$



$$\begin{cases} \mathcal{L} = \varepsilon \Delta - \nabla V \cdot \nabla = \varepsilon e^{V/\varepsilon} \nabla \cdot e^{-V/\varepsilon} \nabla \\ \mathcal{L}^\dagger = \varepsilon \nabla \cdot e^{-V/\varepsilon} \nabla e^{V/\varepsilon} \end{cases}$$

- $\mathcal{L}^\dagger(e^{-V/\varepsilon}) = 0 \Rightarrow \pi(x) = \frac{1}{Z} e^{-V(x)/\varepsilon}$  inv. density
- $\mathcal{L}^\dagger e^{-V/\varepsilon} = e^{-V/\varepsilon} \mathcal{L} \Rightarrow (\pi e^{2t})^\dagger = \pi e^{2t}$  reversibility  
 (detailed balance)

$$\tau := \inf \{ t > 0 : \|x_t - y\| < \delta \} \quad \mathbb{E}^x[\tau] = ?$$

Arrhenius 1889:  $\mathbb{E}^x[\tau] \approx e^{[V(z)-V(x)]/\varepsilon}$  i.e.  $\lim_{\varepsilon \rightarrow 0} \varepsilon \log \mathbb{E}^x[\tau] = V(z) - V(x)$

Eyring 1935, Kramers 1940:  $\mathbb{E}^x[\tau] = \frac{2\pi}{|\lambda_0(z)|} \sqrt{\frac{|\det \text{Hess } V(z)|}{\det \text{Hess } V(x)}} e^{-[V(z)-V(x)]/\varepsilon} [1 + o_\varepsilon(1)]$   
 ↑  
 unique negative ev of Hess V(z)

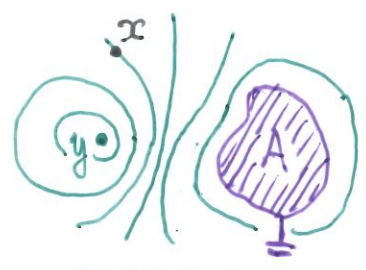
- Proofs:
- Arrhenius: large deviations [Freidlin & Wentzell ~ 1970]
  - Eyring-Kramers:
    - WKB theory
    - potential theory [Baviera, Eckhoff, Gayraud, Klein 2004]
    - Witten Laplacian [Helffer, Klein, Nier 2005]

## Potential theory

$$\tau_A = \inf \{t \geq 0 : x_t \in A\}$$

a)  $W_A(x) = \mathbb{E}^x[\tau_A]$  solves 
$$\begin{cases} -\mathcal{L} W_A(x) = 1 & x \in A^c \\ W_A(x) = 0 & x \in A \end{cases}$$

Green's function: 
$$\begin{cases} -\mathcal{L} G_{A^c}(x, y) = \delta(x-y) & x \in A^c \\ G_{A^c}(x, y) = 0 & x \in A \end{cases}$$



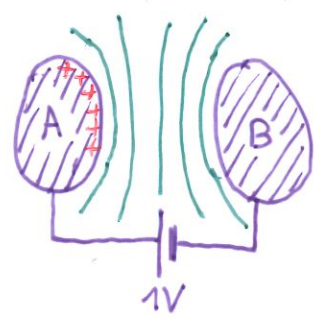
$\epsilon = 1, V = 0$ :  
electric potential

$$\Rightarrow W_A(x) = \int_{A^c} G_{A^c}(x, y) dy$$

Rem: 
$$e^{-V(x)/\epsilon} G_{A^c}(x, y) = e^{-V(y)/\epsilon} G_{A^c}(y, x)$$

b) Comittor:

$h_{AB}(x) = \mathbb{P}^x(\tau_A < \tau_B)$  solves 
$$\begin{cases} \mathcal{L} h_{AB}(x) = 0 & x \in (A \cup B)^c \\ h_{AB}(x) = 1 & x \in A \\ h_{AB}(x) = 0 & x \in B \end{cases}$$



$$\Rightarrow h_{AB}(x) = \int_{\partial A} G_{B^c}(x, y) \underbrace{\rho_{AB}(dy)}_{\text{equil. measure on } \partial A}$$

c) Capacity:  $cap(A, B) = \int_{\partial A} e^{-V(y)/\epsilon} \rho_{AB}(dy) \Rightarrow M_{AB}(dy) = \frac{e^{-V(y)/\epsilon} \rho_{AB}(dy)}{cap(A, B)}$   
 prob. on  $\partial A$

Dirichlet principle:  $cap(A, B) = \epsilon \int_{(A \cup B)^c} \|\nabla h_{AB}(x)\|^2 e^{-V(x)/\epsilon} dx =: \Phi_{(A \cup B)^c}(h_{AB})$

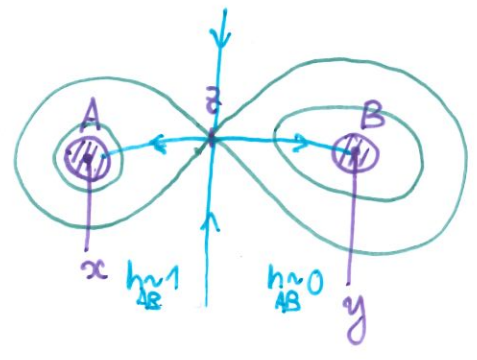
$cap(A, B) = \inf_{\substack{h|_A=1 \\ h|_B=0}} \Phi_{(A \cup B)^c}(h)$   
 $\sim \langle h_{AB}, -\mathcal{L} h_{AB} \rangle$   
 Dirichlet form

d) 
$$\int_{B^c} h_{AB}(x) e^{-V(x)/\epsilon} dx = \int_{B^c} \int_{\partial A} G_{B^c}(x, y) e^{-V(x)/\epsilon} \rho_{AB}(dy) dx$$
  

$$= \int_{\partial A} e^{-V(y)/\epsilon} W_B(y) \rho_{AB}(dy) = cap(A, B) \mathbb{E}^{M_{AB}}[\tau_B]$$

$$\Rightarrow \mathbb{E}^{M_{AB}}[\tau_B] = \frac{1}{cap(A, B)} \int_{B^c} h_{AB}(x) e^{-V(x)/\epsilon} dx$$

Proof of E-K for double-well potential:



- $\text{cap}(A, B) \simeq \varepsilon \frac{|\lambda_0(z)|}{2\pi\varepsilon} \sqrt{\frac{(2\pi\varepsilon)^d}{|\det \text{Hess} V(z)|}} e^{-V(z)/\varepsilon}$
- $\int_{\mathcal{B}^c} h_{AB}(x) e^{-V(x)/\varepsilon} dx \simeq \sqrt{\frac{(2\pi\varepsilon)^d}{|\det \text{Hess} V(x)|}} e^{-V(x)/\varepsilon}$
- $\mathbb{E}^{M_{AB}}[\tau_B] \simeq \mathbb{E}^x[\tau_B]$  (Harnack or coupling)

2. Allen-Cahn in dimension 1

$\partial_t \phi = \Delta \phi + \phi - \phi^3 + \sqrt{2\varepsilon} \xi \leftarrow$  space-time white noise:  
 $\langle \xi, \varphi \rangle$  centred Gaussian r.v.  
 $\mathbb{E}(\langle \xi, \varphi_1 \rangle \langle \xi, \varphi_2 \rangle) = \langle \varphi_1, \varphi_2 \rangle$

$\phi = \phi(t, x) \quad t \geq 0, x \in \mathbb{T} = \mathbb{R}/L\mathbb{Z}$

Potential:  $V(\phi) = \int_0^L \left[ \frac{1}{2} \nabla \phi(x)^2 - \frac{1}{2} \phi(x)^2 + \frac{1}{4} \phi(x)^4 \right] dx$

$\Rightarrow \lim_{h \rightarrow 0} \frac{V(\phi + h\psi) - V(\phi)}{h} = \int_0^L [\nabla \phi(x) \nabla \psi(x) - \phi(x)\psi(x) + \phi(x)^3 \psi(x)] dx$   
 $= - \langle \Delta \phi + \phi - \phi^3, \psi \rangle$

Stationary solutions:  $\phi_0(x) \equiv 0, \phi_{\pm}(x) \equiv \pm 1$ , additionnal non-const. solutions for  $L > 2\pi$

Fourier basis:  $\phi(t, x) = \sum_{k \in \mathbb{Z}} z_k(t) e_k(x) \Rightarrow \boxed{dz_t = -\nabla \hat{V}(z_t) dt + \sqrt{2\varepsilon} dW_t}$

$L < 2\pi$ :  $V =$  double-well potential with saddle in  $\phi_0 = 0$

Hessian at  $\phi_0 = 0$ :  $V(\phi) = \frac{1}{2} \langle \phi, \underbrace{(-\Delta - 1)}_{\text{Hess} V(\phi_0)} \phi \rangle + o(\phi^4)$

$\Rightarrow \mathbb{E}^{\phi_0}[\tau_{\phi_+}] \stackrel{?}{=} \frac{2\pi}{1-1} \underbrace{\sqrt{\frac{|\det(-\Delta-1)|}{|\det(-\Delta+2)|}}}_{\infty/\infty} e^{\frac{[V(\phi_0) - V(\phi_+)]/\varepsilon}{1 + o_\varepsilon(1)}}$   
 compatible with LDP by Farris & Jona-Lasinio

ev of  $\Delta-1$ :  $\lambda_k = (k \frac{2\pi}{L})^2 - 1$        $\Delta_{\perp}$ : projection of  $\Delta$  on  $k \neq 0$   
 $\Delta+2$ :  $\lambda_k + 3$

$$\det([-\Delta_{\perp} - 1][-\Delta_{\perp} + 2]^{-1}) = \det([-\Delta_{\perp} + 2 - 3][-\Delta_{\perp} + 2]^{-1})$$

$$= \det(1 - 3[\Delta_{\perp} + 2]^{-1}) \quad \text{Fredholm determinant}$$

$$\log \det([-\Delta_{\perp} - 1][-\Delta_{\perp} + 2]^{-1}) = \text{Tr} \log(1 - 3[\Delta_{\perp} + 2]^{-1}) = - \sum_{n \geq 1} \frac{3^n}{n} \text{Tr}([-\Delta_{\perp} + 2]^{-n})$$

$$\text{Tr}([-\Delta_{\perp} + 2]^{-n}) = \sum_{k \neq 0} \frac{1}{[(k \frac{2\pi}{L})^2 + 2]^n} \leq \frac{C}{[(\frac{2\pi}{L})^2 + 2]^n}$$

Theorem: [B, Gentz, EJP 2013] E-k formula holds

Proof: uses spectral Galerkin approx.      Rem: prefactor =  $\frac{\sin L}{\sqrt{2} \sinh(\sqrt{2}L)}$  (Euler)

3. Allen-Cahn in dimension 2

$$\left. \begin{aligned} \text{Tr} [(-\Delta_{\perp} + 2)^{-1}] &\sim \sum_{k \neq (0,0)} \frac{1}{\|k\|^2} \sim \int_1^{\infty} \frac{r dr}{r^2} = +\infty \\ \text{Tr} [(-\Delta_{\perp} + 2)^{-2}] &\sim \sum_{k \neq (0,0)} \frac{1}{\|k\|^4} \sim \int_1^{\infty} \frac{r dr}{r^4} < +\infty \end{aligned} \right\} \begin{aligned} &(-\Delta_{\perp} + 2)^{-1} \text{ is} \\ &\text{not trace-class} \\ &\text{but Hilbert-Schmidt} \end{aligned}$$

Theorem: [Da Prato, Debussche 2003]

$$\partial_t \phi = \Delta \phi + \phi - [\phi^3 - 3\epsilon C_N \phi] + \sqrt{2\epsilon} \xi_N$$

with  $C_N \sim \log N$  admits limit as  $N \rightarrow \infty$        $\xi_N$  mollified on scale  $1/N$

$$C_N = \mathbb{E} \|\phi_N\|_{L^2}^2 = \frac{1}{2} \text{Tr} (P_N [-\Delta + 1]^{-1})$$

where  $\phi_N$  is Gaussian free field (GFF):  $\Delta \phi_N - \phi_N + \xi_N = 0$

Renormalized potential:  $V_N(\phi) = \frac{1}{2} \int_{\mathbb{T}^2} [\nabla \phi^2 - \phi^2] dx + \frac{1}{4} \int_{\mathbb{T}^2} [\phi^4 - 6C_N \phi^2 + 3C_N^2] dx$

$\underbrace{\quad}_{=: \phi^4: \text{ (Wick)}}$

$\Rightarrow$  new prefactor:  $\det([-\Delta_{\perp} - 1][-\Delta_{\perp} + 2]^{-1}) e^{3C_N}$   
 $= \det(1 - 3[-\Delta_{\perp} + 2]^{-1}) e^{3\text{Tr}[-\Delta + 2]^{-1}}$       Carleman-Fredholm det

Theorem: [B, Di Gesù, Weber, EJP 2017] unif. upper/lower bds on EK